

Nucleon Magnetic Properties from Lattice QCD with the Background Field Method

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The magnetic moment and magnetic polarisability of the neutron and proton are investigated using the uniform background-field method and lattice QCD. The results are calculated using $32^3 \times 64$ dynamical QCD lattices provided by the PACS-CS collaboration through the ILDG. Methods of isolating the energy of the hadrons accounting for Landau level energies are explored. Effective energy shifts are formed from combinations of correlation functions and their efficiency in isolating the magnetic moment and polarisability terms is determined.

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1. Introduction

The magnetic properties of the nucleon such as the magnetic moment and polarisability can be accessed using lattice QCD with the background field method [1–6]. The magnetic polarisability is a measure of the deformation of a system of charges in an external magnetic field. This deformation causes an energy shift in the particle which can be determined by the energy field relation [1, 7]

$$E(B) = M + \vec{\mu} \cdot \vec{B} + \frac{|qeB|}{2M} - \frac{4\pi}{2} \beta B^2 + \mathcal{O}(B^3). \quad (1.1)$$

Previous studies [5] have faced difficulty in extracting a reliable signal for the polarisability. This is due to the polarisability being a second order effect and the complication of the Landau levels which cannot be easily isolated from the polarisability.

The Landau levels are a series of energy levels arising from a charge or system of charges in an external magnetic field. Hence the charged proton has a Landau level that must be accounted for. In the absence of QCD, the constituent quarks would have individual Landau levels. It is an open question as to the extent to which this effect remains (if it all) in the presence of QCD interactions.

2. Background Field Method

To introduce a background field on the lattice, first consider the continuum case. Here the covariant derivative is modified by the addition of an electromagnetic coupling

$$D_\mu \rightarrow D'_\mu = \partial_\mu + g G_\mu + qe A_\mu, \quad (2.1)$$

Where qe is the charge on the fermion field and A_μ is the electromagnetic four-potential. Discretising this additional term in the same way as the usual gauge fields [8] results in the gauge links being multiplied by an exponential phase factor

$$U_\mu(x) \rightarrow U_\mu(x)^{(B)} = e^{iaqeA_\mu(x)} U_\mu(x). \quad (2.2)$$

Thus far the electromagnetic gauge potential has not been specified uniquely. In order to obtain a magnetic field along the \hat{z} axis, a potential

$$A_x = -B\hat{y}, \quad (2.3)$$

is used over the interior of the $N_x \times N_y \times N_z \times N_t$ lattice. The periodic boundary conditions of the lattice require a non-trivial potential to ensure that the field is uniform over the entirety of the lattice. This requirement produces a quantisation condition on the magnetic field strength [5]

$$|q_d e B| = \frac{2\pi k_d}{N_x N_y a^2}, \quad (2.4)$$

where k_d is an integer governing the field strength.

3. Simulation Details

The calculations detailed here use $2 + 1$ flavour dynamical QCD configurations provided by the PACS-CS collaboration [9] through the International Lattice Data Grid [10]. These lattices have dimensions $32^3 \times 64$ with $\beta = 1.9$ and a physical lattice spacing of $a = 0.0907(13)$ fm. A clover fermion action and Iwasaki gauge action are used. A single value of the light-quark hopping parameter, $k_{ud} = 0.13754$ corresponding to a pion mass of $m_\pi = 413$ MeV is used in this study. The lattice spacing for this mass was set using the Sommer scale with $r_0 = 0.49$ fm. The configuration ensemble size was 450.

To be able to use Eq. (1.1) to extract the polarisabilities, correlation functions at four distinct magnetic field strengths are calculated. As the u and d quarks have different signs, separate propagators at different field strengths must be calculated for each distinct field strength. These correspond to $k_d = 0, \pm 1, \pm 2, \pm 3, \pm 4, \pm 6$ in Eq. (2.4).

The configurations used in this study did not include a background field when generated. Hence the only quarks which feel the presence of the external magnetic field are the valence quarks of the hadrons. To include the background field on the configurations requires separate gauge field configurations for each field strength. This is prohibitively expensive and also destroys the advantageous correlations between the field strengths.

4. Magnetic Polarisability

From correlation functions calculated with a background field in place, the magnetic polarisability can be extracted. To do this consider the energy-field relation in Eq. (1.1). We wish to remove the $\vec{\mu} \cdot \vec{B}$ and M terms. This can be done by using the spin dependence of the $\vec{\mu} \cdot \vec{B}$ term and the zero-field correlator. Taking a combination of spin orientations and field strengths produces the desired result for the energy shift

$$\begin{aligned} \Delta E_p(B) &= \frac{1}{2} (E_\uparrow(B) + E_\downarrow(B) - E_\uparrow(0) - E_\downarrow(0)) \\ &= \frac{|qeB|}{2M} - \frac{4\pi}{2} \beta B^2. \end{aligned} \quad (4.1)$$

A superior method with which to extract this energy is to take a ratio of the correlators directly. This has the advantage of allowing correlated errors to cancel prior to fitting

$$R_p(B, t) = \left(\frac{G_\downarrow(B+, t) + G_\uparrow(B-, t)}{G_\downarrow(0, t) + G_\uparrow(0, t)} \right) \left(\frac{G_\downarrow(B-, t) + G_\uparrow(B+, t)}{G_\downarrow(0, t) + G_\uparrow(0, t)} \right). \quad (4.2)$$

Here the \uparrow and \downarrow represent spin up and down while $B\pm$ represents magnetic fields in the positive and negative \hat{z} directions. From this ratio, an effective energy shift can be extracted in an analogous way to an effective mass.

5. Quark Projection

The $qeB/(2M)$ term in Eq. (4.1) is a Landau level term, it corresponds to the lowest lying Landau level of the hadron. The Landau levels are a superposition of energy levels [11]

$$E^2 = m^2 + |qeB| (2\nu + 1) - q|eB|s + p_z^2,$$

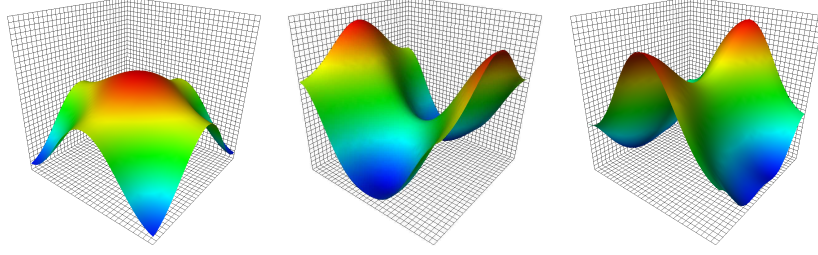


Figure 1: Lowest-lying eigenmode probability densities of the lattice Laplacian operator in a constant background magnetic field oriented in the z direction are plotted as a function of the x, y coordinates. The mode for the lowest quantised magnetic field strength relevant to the d quark (left) is compared with the two degenerate eigenmodes of the second quantised magnetic field strength (middle, right) relevant to the u quark.

caused by the motion of a charged particle in an external magnetic field. Here $v = 0, 1, 2, \dots$, spin parameter $s = \pm 1$ and p_z is the component of momentum in the \hat{z} direction. The charged quarks are also in an external magnetic field and in the absence of QCD would also have Landau level energies. It is possible to obtain these Landau levels using the eigenmodes of the lattice Laplacian operator for each quark in a background magnetic field. A sample of the eigenmodes for the smallest field strength are presented in Figure 1. It is clear from Figure 1 that a particle at the centre of the Lattice will have little overlap with the Landau levels; hence knowledge of the eigenmodes may prove advantageous when constructing quark operators on the lattice.

5.1 Eigenmode Projections

Eigenmodes of the lattice Laplacian operator are calculated where no QCD effects are present; only the QED background field is present.

Once the eigenmodes λ_i have been obtained, the quark propagators can be projected to the eigenmodes at both the source and the sink. Projection operators P_{QED}^n are defined

$$P_{QED}^n(x, y) = \sum_{i=1}^{n=|3q_f k_d|} \langle x | \lambda_i \rangle \langle \lambda_i | y \rangle, \quad (5.1)$$

where q_f is the fractional quark charge. The propagator is then projected at the sink using these projection operators as,

$$S(x, y) = P_{QED}(x, z) S(z, y) \quad (5.2)$$

Any combination of projection operators and smearing can be used at both the source and the sink although only a few have been investigated here.

5.2 Polarisability energy shifts

The energy shift due to the polarisability is smaller than that due to the magnetic moment and it also contains contributions from the Landau levels. These features make the polarisability considerably more challenging to extract than the magnetic moment. In order to extract a polarisability

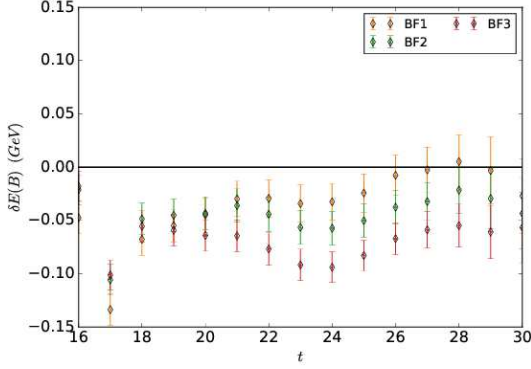


Figure 2: Polarisability energy shift for all field strengths for neutron with smeared source and QED eigenmode projected sink.

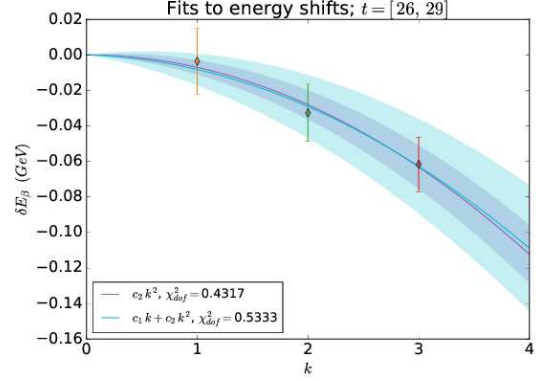


Figure 3: Quadratic and quadratic + linear fits of the energy shift to the field strength for the neutron with smeared source and QED eigenmode projected sink. Here $k = k_d$.

from the energy shifts, a relevant function must be fitted as a function of the field strength - or the field strength quanta k_d . However it is only sensible to do this where acceptable constant fits to the energy shift at each field strength can be adequately performed. The restrictions imposed on fitting are,

1. Constant fits to Eq. (4.1) as a function of t must be acceptable;
2. Relevant fits to Eq. (4.1) as a function of B must be acceptable;
3. Only the same fit window accross all field strengths is considered.

These measures help ensure that the final fits produced are free from bias due to the selection of fit window. Best results for the neutron were found using a spatially smeared source with 100 sweeps of Gaussian smearing and a QED eigenmode projected sink.

For the neutron, it is clear from Figure 3 that a fit which is quadratic-only is sufficient to describe the energy shift as a function of the field strength B . That is, the neutron doesn't have a Landau level energy term. This is as expected as the neutron is a neutrally charged particle. From the fitted curves the effective charge of the hadron and the magnetic polarisability can be extracted. The quadratic only fit results in a value of $\beta_n = 1.31(38) \times 10^{-4} \text{ fm}^3$ for our pion mass of 413 MeV.

6. Magnetic Moment

The magnetic moment of a system of charged particles, $\vec{\mu}$ is related to the tendency of the system to become aligned with the external magnetic field. Returning to Eq. (1.1), the magnetic moment is a first order term in B . This results in a much larger shift in energies than the polarisability term making magnetic moments easier to extract. Of additional use in isolating the magnetic moment energy shift is the spin and field direction dependence.

By forming combinations of spin up and down correlation functions at both positive and negative field strengths the magnetic moment term of the energy shift can be efficiently isolated. This

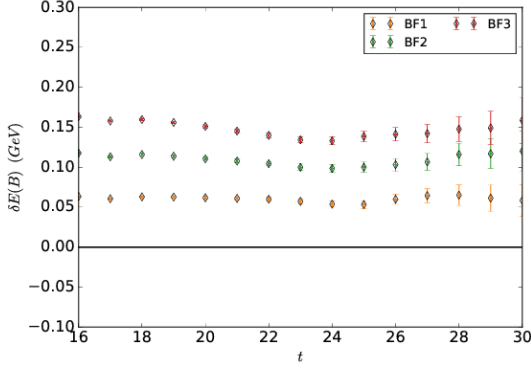


Figure 4: Magnetic moment energy shift for the three field strengths considered for the neutron with a smeared source and a QED-eigenmode projected sink.

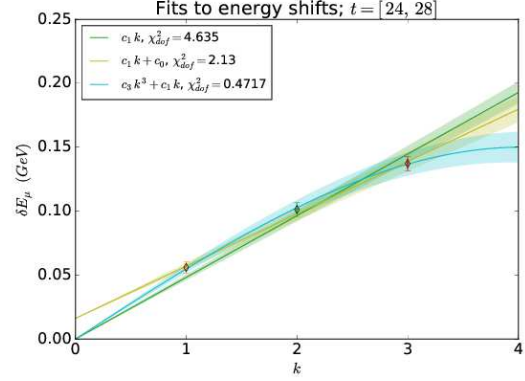


Figure 5: Linear and linear + cubic fits of the energy shift to the field strength for the neutron with a smeared source and a QED-eigenmode projected sink.

can be done using the ratio,

$$\begin{aligned}
 R_m(B, t) &= \left(\frac{G_\downarrow(B+, t) + G_\uparrow(B-, t)}{G_\downarrow(0, t) + G_\uparrow(0, t)} \right) \left(\frac{G_\downarrow(0, t) + G_\uparrow(0, t)}{G_\downarrow(B-, t) + G_\uparrow(B+, t)} \right) \\
 &= \left(\frac{G_\downarrow(B+, t) + G_\uparrow(B-, t)}{G_\downarrow(B-, t) + G_\uparrow(B+, t)} \right)
 \end{aligned} \tag{6.1}$$

In an analogous way to an effective mass, a magnetic moment energy shift can be found,

$$\begin{aligned}
 \Delta E_m(B, t) &= \frac{1}{\delta t} \log \left(\frac{R_\mu(B, t)}{R_\mu(B, t + \delta t)} \right) \\
 &= -\mu B
 \end{aligned} \tag{6.2}$$

This formulation of the energy shift has the advantage that it removes many of the correlated errors between spin orientations.

In the same manner as the magnetic polarisability, magnetic moment energy shifts have constant plateaus fitted to them. This time linear or linear + cubic terms are considered. This cubic term is appropriate as it corresponds to the next lowest order term in Eq. (6.2).

From Figures 5 and 7 it is clear that the cubic term is necessary in order to adequately fit the energy shifts. This suggests that the third field strength in particular is becoming too large for the energy relation in Eq. (1.1) to fully describe the system. This could be remedied by using a larger lattice volume and corresponding smaller field strengths.

The magnetic moments for the proton and neutron linear + cubic fits are shown in Table 1. Good agreement is seen with results from the alternative three-point function method on the same lattices and at the same pion mass [12].

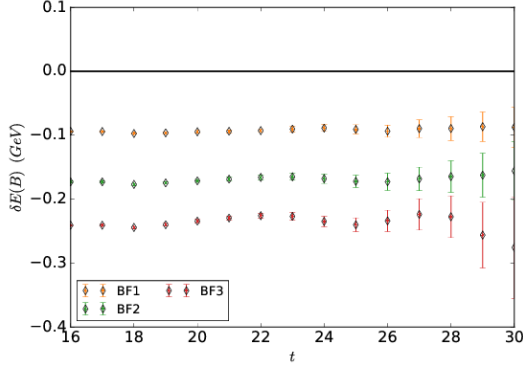


Figure 6: Magnetic moment energy shift for the three field strengths considered for proton with smeared source and QED eigenmode projected sink.

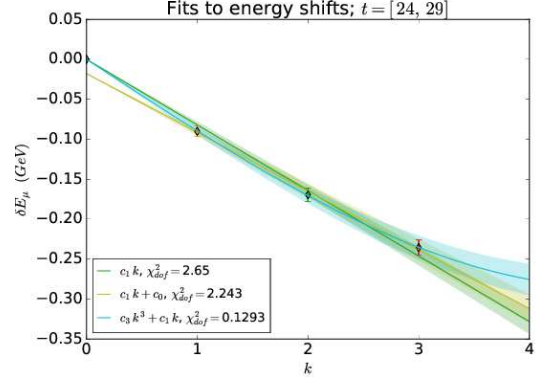


Figure 7: Linear and linear + cubic fits of the energy shift to the field strength for the proton with smeared source and QED eigenmode projected sink.

Table 1: Proton and neutron magnetic moments at a pion mass of 413 MeV

| | This Work | 3PT Method |
|---------|-------------------|--------------------|
| proton | $2.20(16) \mu_N$ | $2.184(22) \mu_N$ |
| neutron | $-1.36(10) \mu_N$ | $-1.371(14) \mu_N$ |

7. Conclusion

Through the use of Landau eigenmode projectors in the sinks of the quark propagators, we have been able to observe plateaus in the correlation functions describing the magnetic polarisability for the first time. We have also examined the utility of a Landau level projector for the proton in its final state. The results are encouraging and the refinement of creation and annihilation operators is in progress. Efforts to expand this method to excited and negative parity states are desirable as well as chiral extrapolations to enable confrontation of experiment. The background field method has been shown again to be a useful tool to access magnetic properties on the lattice.

Acknowledgments

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References

- [1] G. Martinelli, G. Parisi, R. Petronzio and F. Rapuano, *The Proton and Neutron Magnetic Moments in Lattice QCD*, *Phys. Lett.* **B116** (1982) 434–436.

- [2] C. W. Bernard, T. Draper, K. Olynyk and M. Rushton, *Lattice QCD Calculation of Some Baryon Magnetic Moments*, *Phys. Rev. Lett.* **49** (1982) 1076.
- [3] M. Burkardt, D. B. Leinweber and X.-m. Jin, *Background field formalism in quantum systems*, *Phys. Lett.* **B385** (1996) 52–56, [hep-ph/9604450].
- [4] F. X. Lee, L. Zhou, W. Wilcox and J. C. Christensen, *Magnetic polarizability of hadrons from lattice QCD in the background field method*, *Phys. Rev.* **D73** (2006) 034503, [hep-lat/0509065].
- [5] T. Primer, W. Kamleh, D. Leinweber and M. Burkardt, *Magnetic properties of the nucleon in a uniform background field*, *Phys. Rev.* **D89** (2014) 034508, [arXiv:1307.1509].
- [6] NPLQCD collaboration, E. Chang, W. Detmold, K. Orginos, A. Parreno, M. J. Savage, B. C. Tiburzi et al., *Magnetic structure of light nuclei from lattice QCD*, *Phys. Rev.* **D92** (2015) 114502, [arXiv:1506.05518].
- [7] B. C. Tiburzi and S. O. Vayl, *Method to extract charged hadron properties from lattice QCD in magnetic fields*, *Phys. Rev.* **D87** (2013) 054507, [arXiv:1210.4464].
- [8] C. Gattringer and C. B. Lang, *Quantum chromodynamics on the lattice*, *Lect. Notes Phys.* **788** (2010) 1–343.
- [9] PACS-CS collaboration, S. Aoki et al., *2+1 Flavor Lattice QCD toward the Physical Point*, *Phys. Rev.* **D79** (2009) 034503, [arXiv:0807.1661].
- [10] M. G. Beckett, B. Joo, C. M. Maynard, D. Pleiter, O. Tatebe and T. Yoshie, *Building the International Lattice Data Grid*, *Comput. Phys. Commun.* **182** (2011) 1208–1214, [arXiv:0910.1692].
- [11] K. Bhattacharya, *Solution of the Dirac equation in presence of an uniform magnetic field*, arXiv:0705.4275.
- [12] B. J. Owen, *A Variational Approach to Hadron Structure in Lattice QCD*. PhD thesis, The University of Adelaide, 2015.